

# Rethinking the Origin of Neutrino Masses

## The Role of Gravity

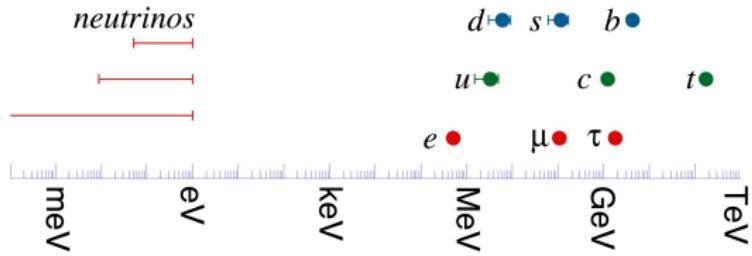
Lena Funcke



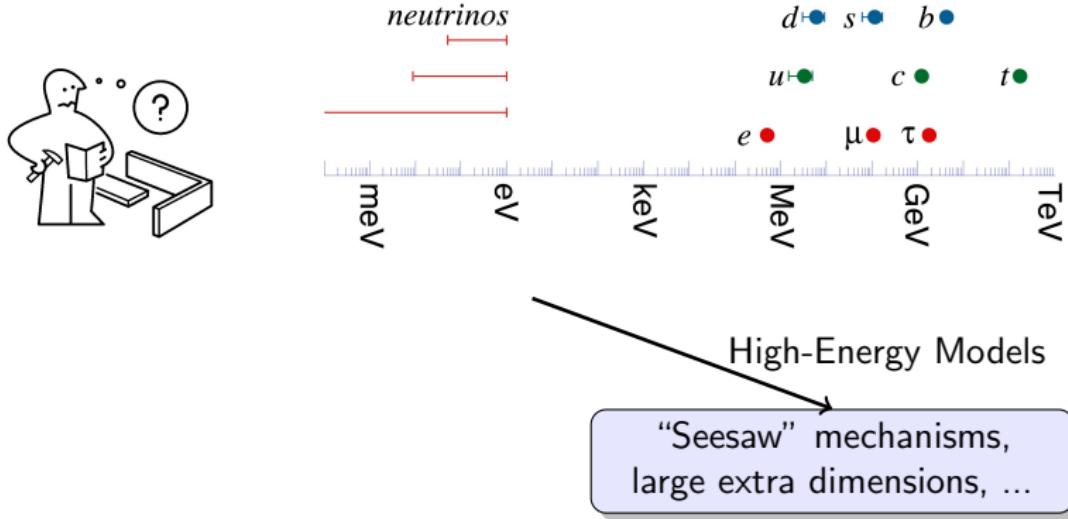
In collaboration with Gia Dvali, Georg Raffelt, and others  
(1602.03191, 1608.08969, 1811.01991, 1905.01264, and ongoing work)

Fermilab Theory Seminar, 3 December 2020

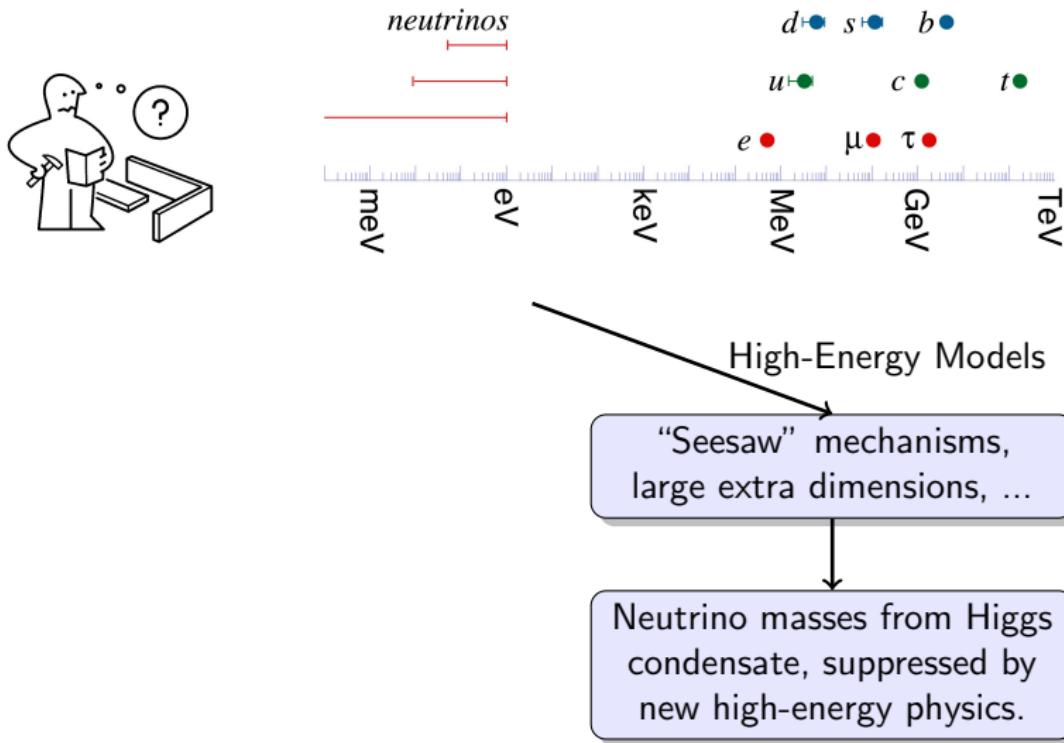
# Question: Origin of Small Neutrino Masses?



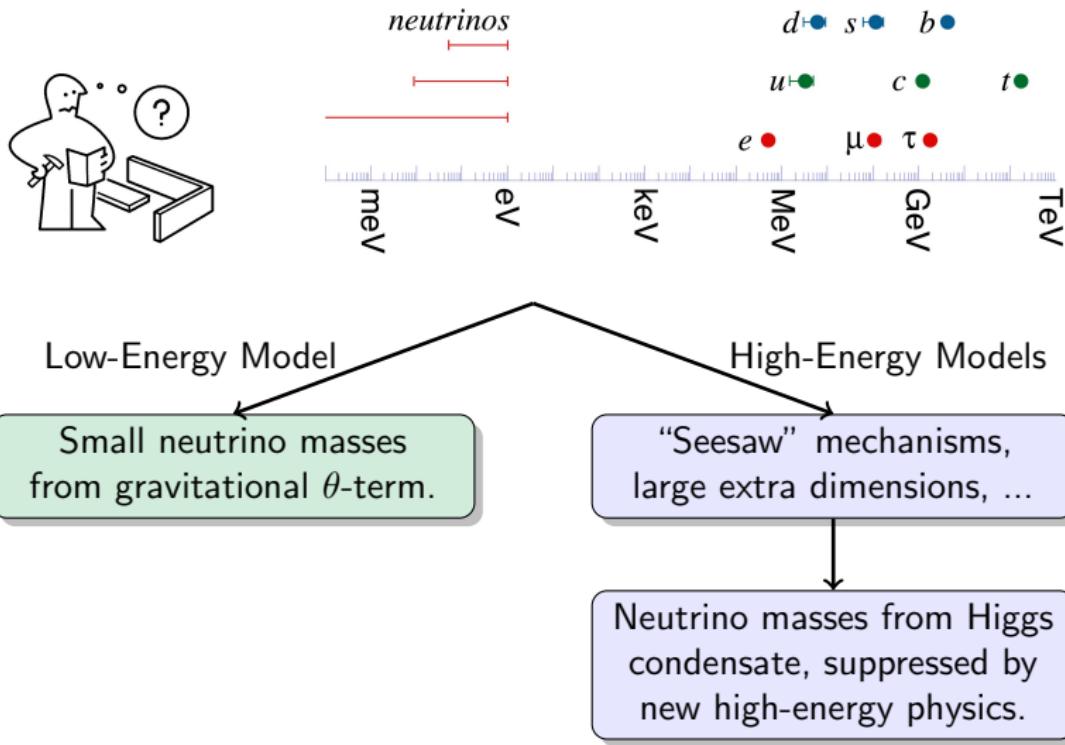
# Question: Origin of Small Neutrino Masses?



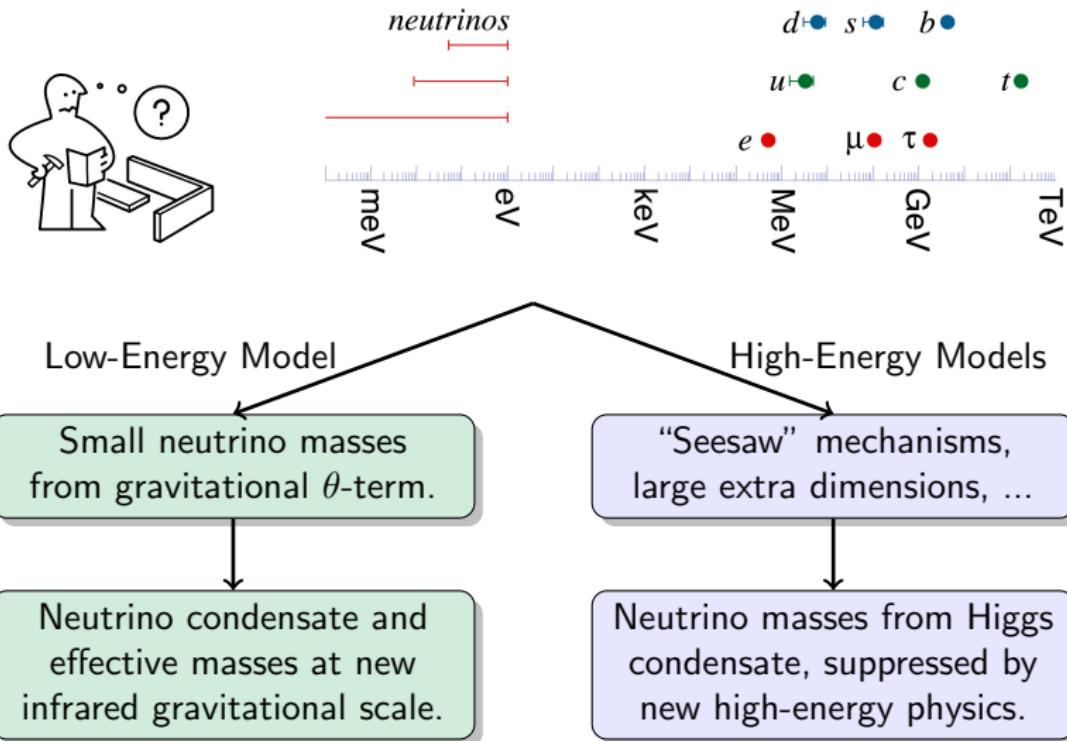
# Question: Origin of Small Neutrino Masses?



# Question: Origin of Small Neutrino Masses?



# Question: Origin of Small Neutrino Masses?



## Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

## Analogy: Non-Perturbative QCD Vacuum



► QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

## Analogy: Non-Perturbative QCD Vacuum



► QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry (Pseudo)Goldstone bosons	$U(3)_V \times U(3)_A \rightarrow U(3)_V$ $1(\eta') + 8(\pi, K, \eta)$

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

## Analogy: Non-Perturbative QCD Vacuum



► QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$	
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$	
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$	
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$	[4]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

[4] Adler (1969); Bell, Jackiw (1969).

## Analogy: Non-Perturbative QCD Vacuum



► QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]
Topological susceptibility	$\langle G \tilde{G}, G \tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q} q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

[4] Adler (1969); Bell, Jackiw (1969). [5] Shifman, Vainshtein, Zakharov (1980).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R \tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]
Topological susceptibility	$\langle G \tilde{G}, G \tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q} q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[4] Adler (1969); Bell, Jackiw (1969). [5] Shifman, Vainshtein, Zakharov (1980).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R \tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]
Topological susceptibility	$\langle G \tilde{G}, G \tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q} q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[4] Adler (1969); Bell, Jackiw (1969). [5] Shifman, Vainshtein, Zakharov (1980).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R \tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]
Topological susceptibility	$\langle G \tilde{G}, G \tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q} q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[4] Adler (1969); Bell, Jackiw (1969). [5] Shifman, Vainshtein, Zakharov (1980).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R \tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$
Fermion flavor symmetry	$U(3)_V \times U(3)_A$ exact if $m_\nu = 0$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]
Topological susceptibility	$\langle G \tilde{G}, G \tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q} q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[4] Adler (1969); Bell, Jackiw (1969). [5] Shifman, Vainshtein, Zakharov (1980).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R\tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$
Fermion flavor symmetry	$U(3)_V \times U(3)_A$ exact if $m_\nu = 0$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = R\tilde{R} + m_\nu \bar{\nu} \gamma_5 \nu$ [6]
Topological susceptibility	$\langle G\tilde{G}, G\tilde{G} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{q}q \rangle \neq 0$ [5]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[5] Shifman, Vainshtein, Zakharov (1980). [6] Delbourgo, Salam (1972); Eguchi, Freund (1976).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R\tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$
Fermion flavor symmetry	$U(3)_V \times U(3)_A$ exact if $m_\nu = 0$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = R\tilde{R} + m_\nu \bar{\nu} \gamma_5 \nu$ [6]
Topological susceptibility	$\langle R\tilde{R}, R\tilde{R} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{\nu} \nu \rangle \neq 0$ [7]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[6] Delbourgo, Salam (1972); Eguchi, Freund (1976). [7] Dvali, LF (2016a).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R \tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$	
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(1)^3$	[7]
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$	
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = R \tilde{R} + m_\nu \bar{\nu} \gamma_5 \nu$	[6]
Topological susceptibility	$\langle R \tilde{R}, R \tilde{R} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{\nu} \nu \rangle \neq 0$	[7]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[6] Delbourgo, Salam (1972); Eguchi, Freund (1976). [7] Dvali, LF (2016a).

# Analogy: Non-Perturbative QCD Vacuum



- QCD:  $\theta$ -term  $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$  is made physical by non-perturbative effects [1].
- Gravity:  $\theta$ -term  $\mathcal{L}_G \supset \theta_G R\tilde{R}$  exists [2], physicality is unknown  $\rightarrow$  assumption!
- Similarities: axial anomalies, topology structures, massive pseudoscalars [3].

Quantity	Gravity with $3\nu$	
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(1)^3$	[7]
(Pseudo)Goldstone bosons	$1(\eta_\nu) + 14(\phi_k)$	[3,7]
Axial $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = R\tilde{R} + m_\nu \bar{\nu} \gamma_5 \nu$	[6]
Topological susceptibility	$\langle R\tilde{R}, R\tilde{R} \rangle_{p \rightarrow 0} \neq 0 \Rightarrow \langle \bar{\nu} \nu \rangle \neq 0$	[7]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979). [2] Deser, Duff, Isham (1980).

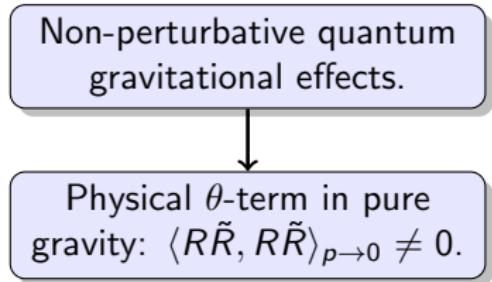
[3] Dvali (2005); Dvali, Jackiw, Pi (2006); Dvali, Folkerts, Franca (2014).

[6] Delbourgo, Salam (1972); Eguchi, Freund (1976). [7] Dvali, LF (2016a).

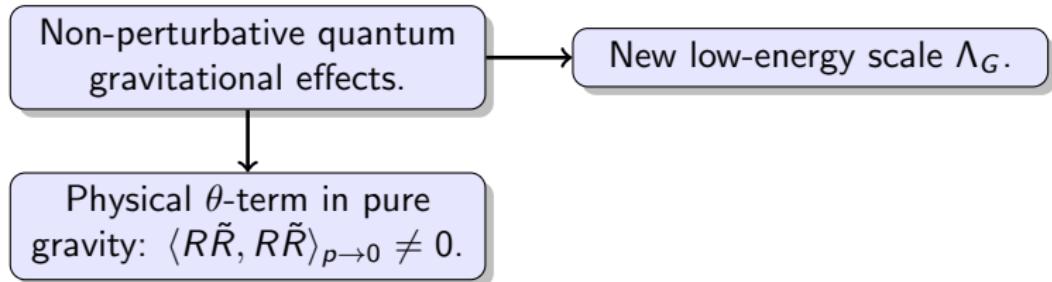
# The Model: Neutrino Condensation

Non-perturbative quantum  
gravitational effects.

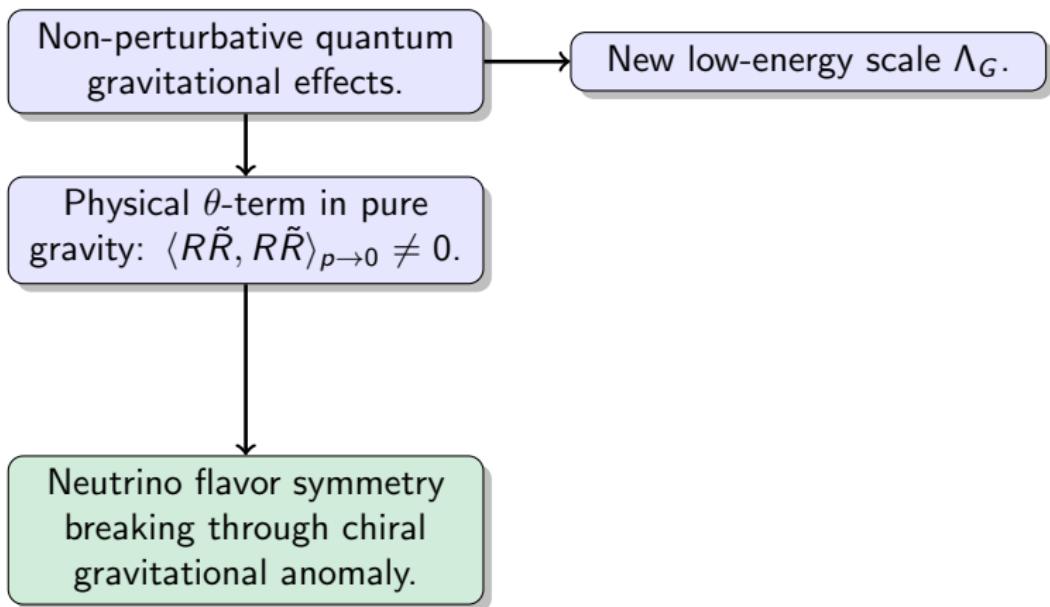
# The Model: Neutrino Condensation



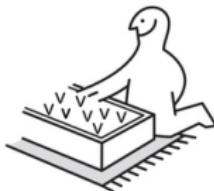
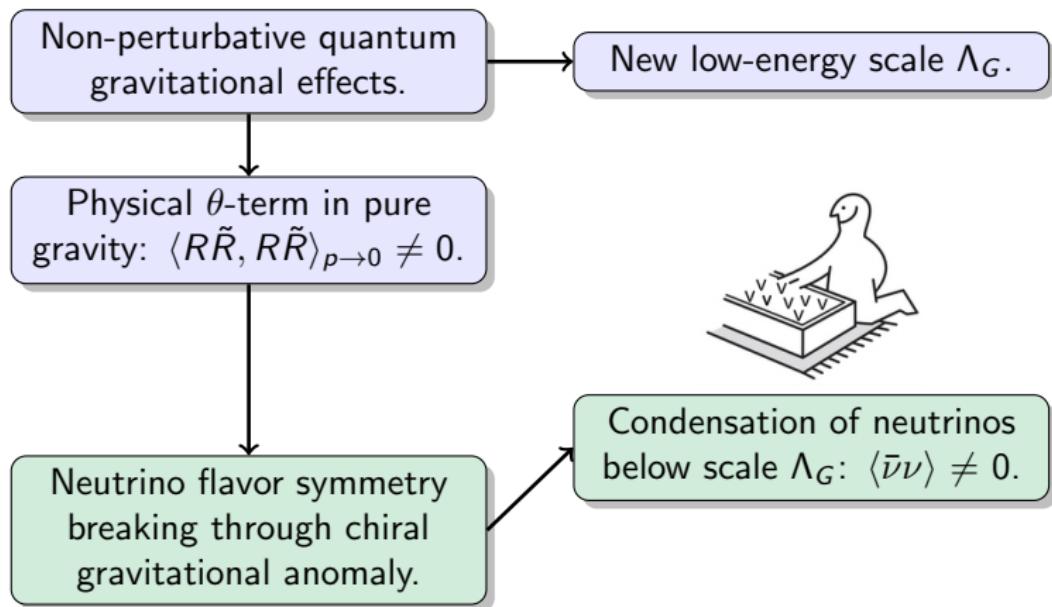
# The Model: Neutrino Condensation



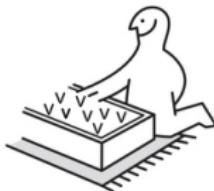
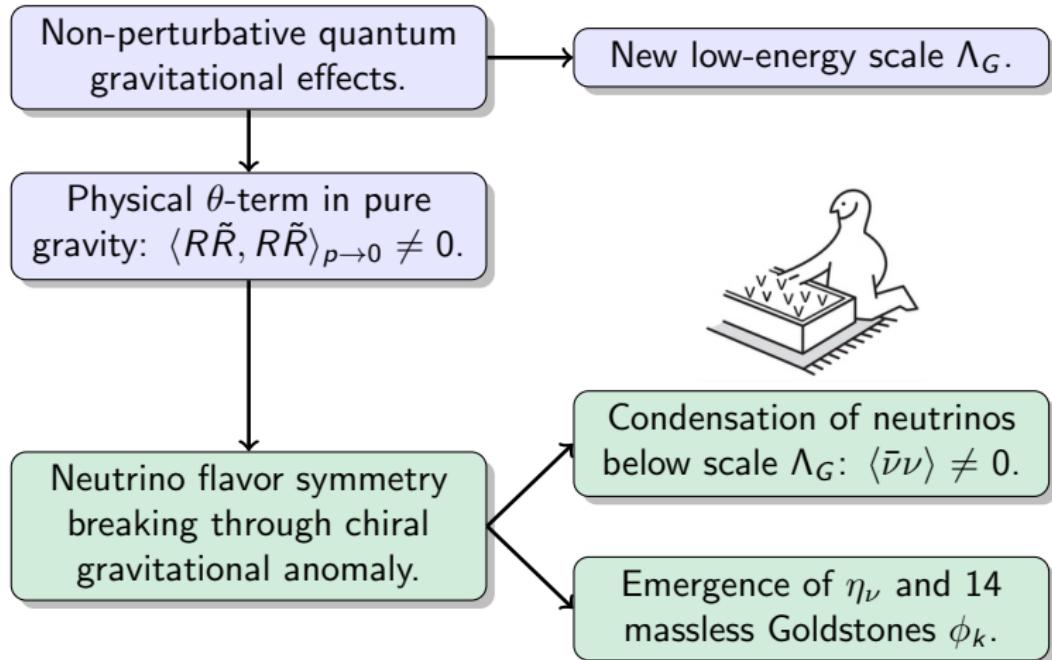
# The Model: Neutrino Condensation



# The Model: Neutrino Condensation

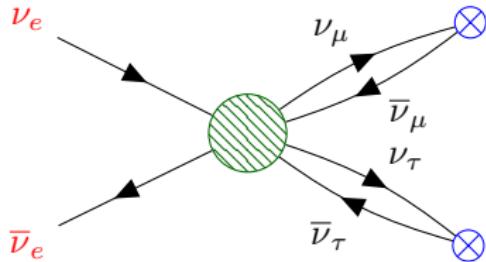


# The Model: Neutrino Condensation



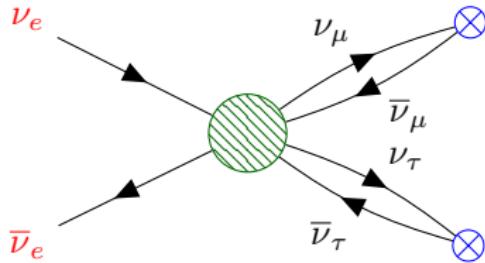
# The Model: Neutrino Mass Generation

- Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.



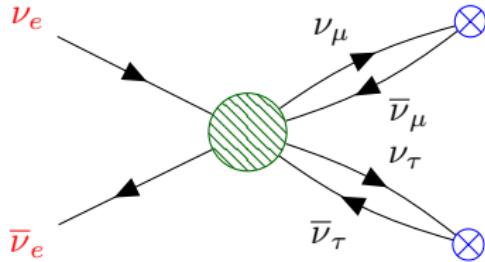
# The Model: Neutrino Mass Generation

- ▶ Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [8].



# The Model: Neutrino Mass Generation

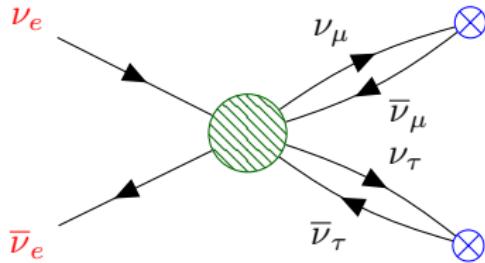
- ▶ Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [8].



- ▶ Effective potential allows for neutrino mass hierarchy.

# The Model: Neutrino Mass Generation

- ▶ Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [8].



- ▶ Effective potential allows for neutrino mass hierarchy.
- ▶ Mechanism works for Dirac and Majorana masses.

[8] 't Hooft (1986).

## Constraints: Symmetry Breaking Scale $\Lambda_G$

Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$

## Constraints: Symmetry Breaking Scale $\Lambda_G$

Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$

$\Lambda_G$



## Constraints: Symmetry Breaking Scale $\Lambda_G$

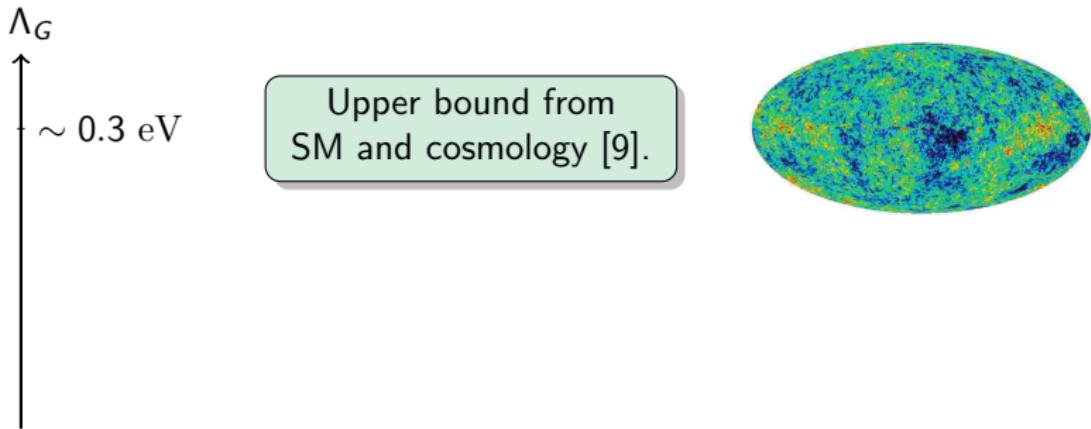
Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$

$\Lambda_G$



# Constraints: Symmetry Breaking Scale $\Lambda_G$

Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$

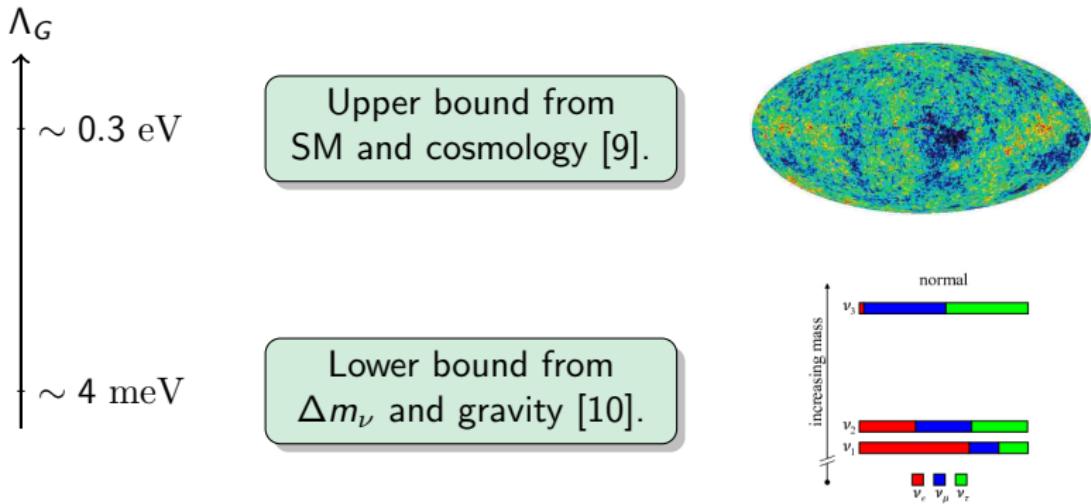


[9] Archidiacono, Hannestad (2014).

Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>]

# Constraints: Symmetry Breaking Scale $\Lambda_G$

Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$

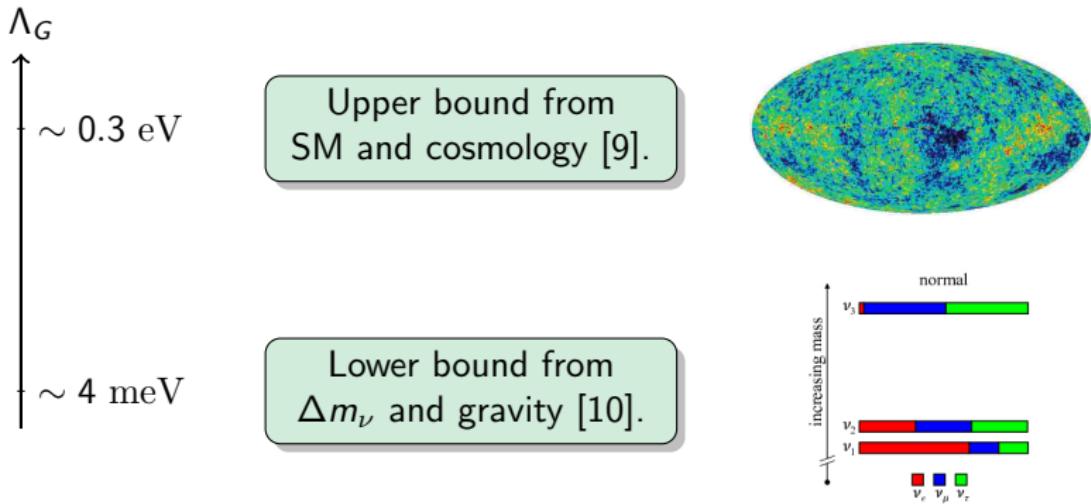


[9] Archidiacono, Hannestad (2014). [10] Tanabashi *et al.* (Particle Data Group) (2018).

Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>] and Patterson (2005).

# Constraints: Symmetry Breaking Scale $\Lambda_G$

Neutrino condensate  $|\langle \bar{\nu} \nu \rangle| \sim \text{scale } \Lambda_G^3 \sim \text{temperature } T_{\chi \text{SB}}^3$



→ Neutrino vacuum condensate  $\langle \bar{\nu} \nu \rangle$  on dark energy scale

[9] Archidiacono, Hannestad (2014). [10] Tanabashi *et al.* (Particle Data Group) (2018).

Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>] and Patterson (2005).

## Phenomenological Implications

Weakened cosmological neutrino mass bounds.

- Relic neutrinos massless until late phase transition at  $T_{\chi\text{SB}} \lesssim \Lambda_G$ .

## Phenomenological Implications

### Weakened cosmological neutrino mass bounds.

- ▶ Relic neutrinos massless until late phase transition at  $T_{\chi\text{SB}} \lesssim \Lambda_G$ .
- ▶ Neutrinos decay & (partially) annihilate  $\rightarrow \sum_i m_{\nu_i} \not\gtrsim 0.12 \text{ eV}$  [11].

# Phenomenological Implications

## Weakened cosmological neutrino mass bounds.

- Relic neutrinos massless until late phase transition at  $T_{\chi\text{SB}} \lesssim \Lambda_G$ .
- Neutrinos decay & (partially) annihilate  $\rightarrow \sum_i m_{\nu_i} \not\gtrsim 0.12 \text{ eV}$  [11].  
 $\Rightarrow$  Masses  $m_{\nu_e} \lesssim 1.1 \text{ eV}$  [12] still allowed, measurable at



[11] Aghanim *et al.* (Planck) (2018). [12] Aker *et al.* (KATRIN) (2019).

Image credit: KATRIN [<http://www.ikp.kit.edu/>].

# Phenomenological Implications

## Weakened cosmological neutrino mass bounds.

- Relic neutrinos massless until late phase transition at  $T_{\chi \text{SB}} \lesssim \Lambda_G$ .
- Neutrinos decay & (partially) annihilate  $\rightarrow \sum_i m_{\nu_i} \not\gtrsim 0.12 \text{ eV}$  [11].
  - $\Rightarrow$  Masses  $m_{\nu_e} \lesssim 1.1 \text{ eV}$  [12] still allowed, measurable at KATRIN.
  - $\Rightarrow$  Light sterile neutrinos cosmologically allowed.



[11] Aghanim *et al.* (Planck) (2018). [12] Aker *et al.* (KATRIN) (2019).

Image credit: KATRIN [<http://www.ikp.kit.edu/>].

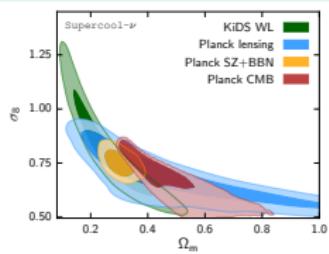
# Phenomenological Implications

## Weakened cosmological neutrino mass bounds.

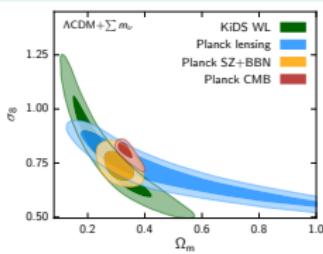
- Relic neutrinos massless until late phase transition at  $T_{\chi \text{SB}} \lesssim \Lambda_G$ .
- Neutrinos decay & (partially) annihilate  $\rightarrow \sum_i m_{\nu_i} \not\gtrsim 0.12 \text{ eV}$  [11].  
 $\Rightarrow$  Masses  $m_{\nu_e} \lesssim 1.1 \text{ eV}$  [12] still allowed, measurable at KATRIN.  
 $\Rightarrow$  Light sterile neutrinos cosmologically allowed.



## Impact on other cosmic parameters.



$\sigma_8(\Omega_m)$  for late  $m_\nu$



vs.  $\sigma_8(\Omega_m)$  for  $\Lambda$ CDM

[11] Aghanim *et al.* (Planck) (2018). [12] Aker *et al.* (KATRIN) (2019).

Image credit: KATRIN [<http://www.ikp.kit.edu/>]. Plots: Lorenz, LF, Calabrese, Hannestad (2018).

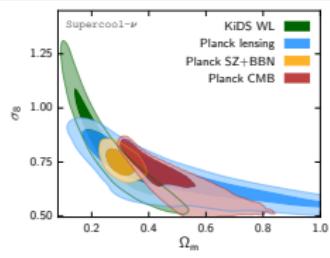
# Phenomenological Implications

## Weakened cosmological neutrino mass bounds.

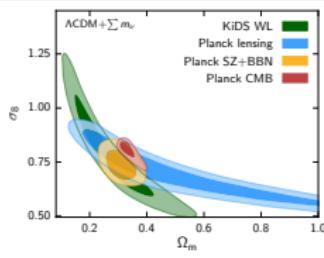
- Relic neutrinos massless until late phase transition at  $T_{\chi \text{SB}} \lesssim \Lambda_G$ .
- Neutrinos decay & (partially) annihilate  $\rightarrow \sum_i m_{\nu_i} \not\gtrsim 0.12 \text{ eV}$  [11].  
 $\Rightarrow$  Masses  $m_{\nu_e} \lesssim 1.1 \text{ eV}$  [12] still allowed, measurable at KATRIN.  
 $\Rightarrow$  Light sterile neutrinos cosmologically allowed.



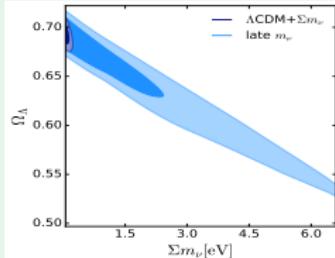
## Impact on other cosmic parameters. Decaying dark energy?



$\sigma_8(\Omega_m)$  for late  $m_\nu$



vs.  $\sigma_8(\Omega_m)$  for  $\Lambda$ CDM



$\Omega_\Lambda(m_\nu)$  for both models

[11] Aghanim *et al.* (Planck) (2018). [12] Aker *et al.* (KATRIN) (2019).

Image credit: KATRIN [<http://www.ikp.kit.edu/>]. Plots: Lorenz, LF, Calabrese, Hannestad (2018).

## Phenomenological Implications

### Astrophysical neutrinos:

- Enhanced neutrino decays: distinct flavor patterns at Earth.

# Phenomenological Implications

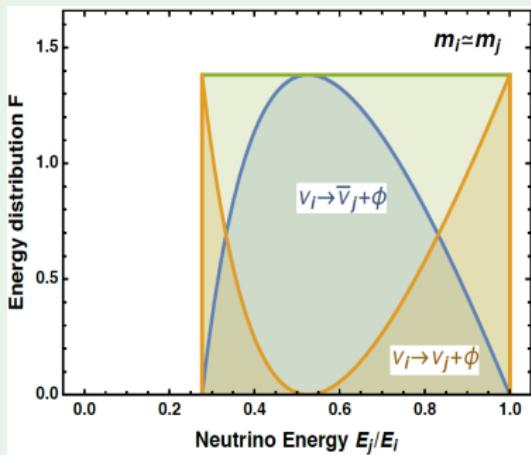
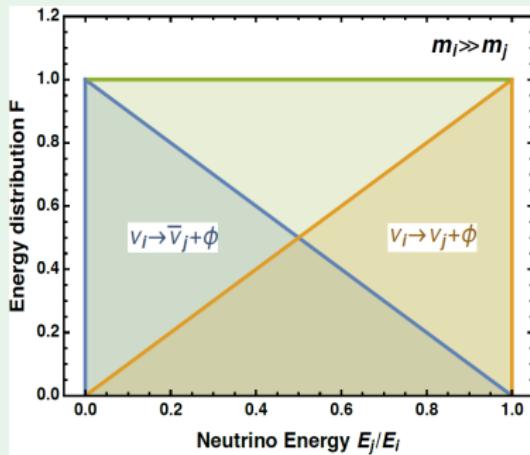
## Astrophysical neutrinos:

- ▶ Enhanced neutrino decays: distinct flavor patterns at Earth.
- ▶ Majorana vs. Dirac neutrinos: different decay channels  $\nu_i \rightarrow \nu_j + \phi$  and  $\nu_i \rightarrow \bar{\nu}_j + \phi$  observable in solar (and future IceCube) data [15].

# Phenomenological Implications

## Astrophysical neutrinos:

- Enhanced neutrino decays: distinct flavor patterns at Earth.
- Majorana vs. Dirac neutrinos: different decay channels  $\nu_i \rightarrow \nu_j + \phi$  and  $\nu_i \rightarrow \bar{\nu}_j + \phi$  observable in solar (and future IceCube) data [15].



# Phenomenological Implications

Gravity and new particle detection:

- ▶ Different polarization intensities of gravitational waves [16].

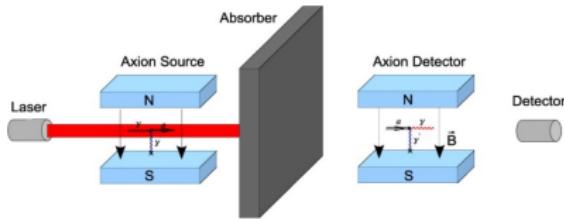


[16] Jackiw, Pi (2003).

# Phenomenological Implications

## Gravity and new particle detection:

- ▶ Different polarization intensities of gravitational waves [16].
- ▶ Searching for new  $\phi$  bosons in axion-like experiments [17].

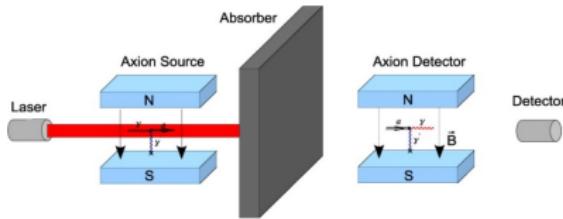


[16] Jackiw, Pi (2003). [17] Dvali, LF, "Domestic Axion" (2016b). Image credits: SXS project and Kim, Carosi (2008).

# Phenomenological Implications

## Gravity and new particle detection:

- ▶ Different polarization intensities of gravitational waves [16].
- ▶ Searching for new  $\phi$  bosons in axion-like experiments [17].



## Ongoing projects:

- ▶ Soft topological defects from neutrino phase transition [18].

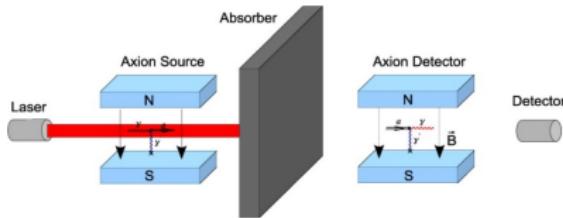
[16] Jackiw, Pi (2003). [17] Dvali, LF, "Domestic Axion" (2016b). Image credits: SXS project and Kim, Carosi (2008).

[18] Dvali, LF, Vachaspati, to appear.

# Phenomenological Implications

## Gravity and new particle detection:

- ▶ Different polarization intensities of gravitational waves [16].
- ▶ Searching for new  $\phi$  bosons in axion-like experiments [17].



## Ongoing projects:

- ▶ Soft topological defects from neutrino phase transition [18].
- ▶ Neutrino mass reconstruction as a function of redshift [19].

[16] Jackiw, Pi (2003). [17] Dvali, LF, "Domestic Axion" (2016b). Image credits: SXS project and Kim, Carosi (2008).

[18] Dvali, LF, Vachaspati, to appear. [19] Lorenz, LF, Calabrese, Löffler, to appear.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- Neutrino condensation.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.



Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.



Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.



Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.
- ▶ More details on arXiv:  
1811.01991 & 1905.01264

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.
- ▶ More details on arXiv:  
1811.01991 & 1905.01264



Thanks for listening!

# Summary

Assumption: pure gravity contains physical  $\theta$ -term.

Theoretical consequences:

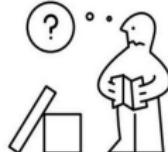
- ▶ Neutrino condensation.
- ▶ Small effective neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ More details on arXiv:  
1602.03191 & 1608.08969

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.
- ▶ More details on arXiv:  
1811.01991 & 1905.01264



Thanks for listening!



Do you have any questions?